1 Sturm-Liouville and Separation of Variables

Recall from Wed June 1st (Lec 13), our original hyperbolic and parabolic equations were

$$w(x) \frac{\partial^2 u}{\partial t^2} + \mathscr{L}[u] = 0, \quad \text{hyperbolic}$$

 $w(x) \frac{\partial u}{\partial t} + \mathscr{L}[u] = 0, \quad \text{parabolic}$

with homogeneous BCs:

$$\alpha_1 u(0,t) - \beta_1 u_x(0,t) = 0, \qquad \alpha_2 u(l,t) + \beta_2 u_x(l,t) = 0$$

and ICs u(x,0) = f(x) and $u_t(x,0) = g(x)$ when appropriate. We then separated the PDEs into the ODEs

$$\frac{1}{w(x)}\mathscr{L}[X](x) = \lambda_k X(x)$$

and

$$T''(t) + \lambda_k T(t) = 0$$
, hyperbolic
 $T'(t) + \lambda_k T(t) = 0$, parabolic

Now the λ_k is the same between the two, so we know that the λ_k s in this second set of ODEs are non-negative and countably-infinite.

For the hyperbolic and parabolic case our final solution will be of the form,

$$u(x,t) = \sum_{k=0}^{\infty} X_k(x) T_k(t)$$

where the X_k s are normalized eigenfunctions via the Sturm-Liouville problem, and the T_k s are from solving the second set of ODEs on the eigenvalues. As long as we can differentiate termwise, this satisfies the original PDE.

For the ICs, if we enforce that $T_k(0) = a_k = (f, X_k)$ as our IC for the separated ODE, we can take $u(x, 0) = \sum_{k=0}^{\infty} a_k X_k(x) = f(x)$ to satisfy the IC for the PDE since the X_k s are complete. If we additionally require in the hyperbolic case that $T'_k(0) = b_k = (g, X_k)$, then we have that $u_t(x, 0) = g(x)$. Hence the ICs are satisfied by solving an IVP in T.

For the BCs note that $u_x = \sum_{k=0}^{\infty} X'_k(x) T_k(t)$ hence

$$\alpha_1 u(0,t) - \beta_1 u_x(0,t) = \sum_{k=0^\infty} (\alpha_1 X_k(0) - \beta_1 X'_k(0)) T_k(t) = 0$$

due to the BC in the Sturm-Liouville problem. For the elliptic case we had

$$w(x)\frac{\partial^2 u}{\partial y^2} - \mathscr{L}[u] = 0$$

Which after taking u(x, y) = X(x) Y(y) gave the SL problem and

$$Y''(y) - \lambda_k Y(y) = 0$$

as ODEs. If we consider the rectangular domain $V = [0 \dots l_x] \times [0 \dots l_y]$, with BCs

$$\alpha_1 u(0, y) - \beta_1 u_x(0, y) = 0, \ \alpha_2 u(l_x, y) + \beta_2 u_x(l_x, y) = 0, \quad u(x, 0) = f(x), \ u(x, l_y) = g(x)$$

then the solution for X_k via the SL problem takes care of the x border's BCs. For the y borders, we have

$$u(x,0) = \sum_{k=0}^{\infty} X_k(x) Y_k(0)$$

so if $Y_k(0) = (f, X_k)$ then the y = 0 condition is satisfied. For $y = l_y$, we require that $Y_k(l_y) = (g, X_k)$. So to satisfy the y BC we must solve the BVP in Y. We should ensure this is always possible: if $\lambda_k \neq 0$, then

$$Y'' - \lambda_k Y \implies Y(y) = c_1 \exp(\sqrt{\lambda_k} y) + c_2 \exp(-\sqrt{\lambda_k} y)$$

Hence

$$\begin{bmatrix} Y(0) \\ Y(l) \end{bmatrix} = \begin{bmatrix} (f, X_k) \\ (g, X_k) \end{bmatrix} = \begin{bmatrix} 1 & 1 \\ \exp(\sqrt{\lambda_k} l_y) & \exp(-\sqrt{\lambda_k} l_y) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \end{bmatrix}$$

where the matrix is invertible for non-zero eigenvalues. If $\lambda = 0$, then the solution is

$$Y'' = 0 \implies Y = a x + b$$

hence by taking Y to be the straight line between (f, X_k) at y = 0 and (g, X_k) at $y = l_y$, the BVP is solved.

2 Examples

In general, we want to find λ -u pairs such that

$$\frac{1}{w}\mathscr{L}[u] = \lambda \, u, \quad \alpha_1 \, u(0) - \beta_1 \, u'(0) = 0, \quad \alpha_2 \, u(l) + \beta_2 \, u'(l) = 0.$$

The general process we will follow is to find v_1 and v_2 such that

$$\frac{1}{w}\mathscr{L}[v_i] = \lambda v_i, \quad v_1(0) = 1, \ v_1'(0) = 0, \quad v_2(0) = 0, \ v_2'(0) = 1$$

If we take $u(x) = \beta_1 v_1(x) + \alpha_1 v_2(x)$, then $u(0) = \beta_1$ and $u'(0) = \alpha_1$. Hence $\alpha_1 u(0) - \beta_1 u'(0) = \alpha_1 \beta_1 - \beta_1 \alpha_1 = 0$. For the second BC we need to use the values of λ to ensure it is satisfied. Our theory above tells us we will always be able to do this (and in so doing will find a countable infinite number of λ s that make this possible). This is more clearly illustrated with an example.